

Problem 3.3: Position and momentum operators

In one-dimensional space the coordinate operator \hat{x} and the momentum operator \hat{p} are defined as

$$\hat{x}\psi(x) = x\psi(x), \quad \hat{p}\psi(x) = -i\hbar\frac{\partial}{\partial x}\psi(x),$$

where $\psi(x)$ is a wave function. Calculate the following commutators:

a) (4 points) $[\hat{p}^2, \hat{x}^2]$ and $[\hat{p}^2, \hat{x}^{-1}]$,

$$[\hat{A}, \hat{B}] = \hat{A}\hat{B} - \hat{B}\hat{A}$$

$$\begin{aligned} [\hat{p}^2, \hat{x}^2]\psi &= (\hat{p}^2 \hat{x}^2 \psi - \hat{x}^2 \hat{p}^2 \psi) = \\ &= -\hbar^2 \frac{\partial^2}{\partial x^2} (x^2 \psi) - x^2 \cdot (-1) \cdot \hbar^2 \frac{\partial^2}{\partial x^2} \psi(x) = \\ &= \hbar^2 \left(-\frac{\partial^2}{\partial x^2} (x^2 \psi) + x^2 \frac{\partial^2}{\partial x^2} \psi(x) \right) = \\ &= \hbar^2 \left(-\frac{\partial}{\partial x} (2x \cdot \psi + x^2 \cdot \frac{\partial}{\partial x} \psi) + x^2 \frac{\partial^2}{\partial x^2} \psi \right) \\ &= \hbar^2 \left(-\left(2\psi + 2x \frac{\partial \psi}{\partial x} + 2x \frac{\partial}{\partial x} \psi + x^2 \frac{\partial^2 \psi}{\partial x^2} \right) + x^2 \frac{\partial^2}{\partial x^2} \psi \right) \\ &= \hbar^2 \left(-2\psi - 4x \frac{\partial \psi}{\partial x} \right) = -2 \cdot \hbar^2 \left(\psi + 2x \frac{\partial \psi}{\partial x} \right) = \\ &= \left(-2 \hbar^2 \left(1 + 2x \frac{\partial}{\partial x} \right) \right) \psi \end{aligned}$$

$$[\hat{p}^2, \hat{x}^2] = -2\hbar^2 \left(1 + 2x \frac{\partial}{\partial x} \right)$$

$$(f \cdot g)' = f'g + f \cdot g'$$

Problem 11: Consider wave packet of a free particle
 The solution of the wave packet is determined by the one-dimensional Schrödinger equation for a free particle.
 In Fourier space the free Schrödinger equation reads

Freie Schrödinger-gleichung
 $i\hbar \frac{\partial}{\partial t} \psi(r, t) = -\frac{\hbar^2}{2m} \Delta \psi(r, t)$

$\int a \cdot f(x) dx = a \cdot \int f(x) dx$

$i \cdot \frac{\partial}{\partial t} \Psi(x, t) = -\frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} \Psi(x, t)$

$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \tilde{\Psi}(k, t) \cdot e^{ikx} dk$

$\frac{1}{\sqrt{2\pi}} i \cdot \frac{\partial}{\partial t} \int_{-\infty}^{\infty} \tilde{\Psi}(k, t) e^{ikx} dk = -\frac{1}{\sqrt{2\pi}} \frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} \int_{-\infty}^{\infty} \tilde{\Psi}(k, t) e^{ikx} dk$

$i \cdot \int_{-\infty}^{\infty} \frac{\partial}{\partial t} \tilde{\Psi}(k, t) \cdot e^{ikx} dk = -\frac{\hbar}{2m} \int_{-\infty}^{\infty} \tilde{\Psi}(k, t) \frac{\partial^2 e^{ikx}}{\partial x^2} dk$

$(e^{ikx})' = e^{ikx} \cdot ik$
 $(e^{ikx})'' = -k^2 e^{ikx}$

$\int_{-\infty}^{\infty} e^{ikx} \cdot \left(i \frac{\partial}{\partial t} \tilde{\Psi}(k, t) - \frac{\hbar}{2m} \tilde{\Psi}(k, t) k^2 \right) dk = 0$

$i \frac{\partial}{\partial t} \tilde{\Psi}(k, t) - \frac{\hbar}{2m} k^2 \tilde{\Psi}(k, t) = 0$

$\frac{\partial}{\partial t} \tilde{\Psi} = \frac{\hbar}{2mi} k^2 \tilde{\Psi}$

$\tilde{\Psi}(k, t) = C \cdot e^{\frac{\hbar}{2mi} k^2 t}$

$\tilde{\Psi}(k, t) = \frac{1}{\sqrt{4\pi}} \cdot e^{-\frac{(k_0+k)^2}{2a}} \cdot e^{\frac{\hbar}{2mi} k^2 t}$

$f'(x) = a \cdot f(x), f(0) = 100$
 $f(x) = 100 \cdot e^{ax}$
 $f(0) = C \cdot e^{a \cdot 0} = C \cdot 1 = C = 100$

Probe: $f'(x) = C \cdot e^{ax} \cdot a = a \cdot C \cdot e^{ax} = a \cdot f(x) \checkmark$

$f'(x) = 2 \cdot f(x), f(0) = 3$
 $f(x) = \frac{2}{1} f(x)$
 $f(x) = C \cdot e^{\frac{2}{1} x} = 3 \cdot e^{\frac{2}{1} x}$
 $f(0) = C \cdot e^{\frac{2}{1} \cdot 0} = C \cdot 1 = 3$
 Probe: $f'(x) = C \cdot \frac{2}{1} \cdot e^{\frac{2}{1} x} = \frac{2}{1} f(x) \checkmark$

c) Show that the time-dependent wave packet $\psi(x, t)$ has the expression (2 points)

$\psi(x, t) = \frac{1}{\sqrt{a\sqrt{\pi}(1+it/\tau)}} \exp\left(-\frac{(x-v_0 t)^2}{2a^2(1+it/\tau)} + i(k_0 x - \omega_0 t)\right)$

where

$\omega_0 = \frac{\hbar k_0^2}{2m}, v_0 = \frac{\hbar k_0}{m}, \tau = \frac{m a^2}{\hbar}$

$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \cdot \int_{-\infty}^{\infty} \tilde{\Psi}(k, t) \cdot e^{ikx} dk =$
 $= \frac{1}{\sqrt{2\pi}} \cdot \frac{1}{\sqrt{4\pi}} \int_{-\infty}^{\infty} e^{-\frac{(k_0+k)^2}{2a}} + \frac{\hbar}{2mi} k^2 t + ikx} dk$

$\int e^x dx = e^x$
 $\int e^{-k^2} dk = ?$

quadratische Ergänzung und Gauß-Integral
 $\int_{-\infty}^{\infty} e^{-\delta(x^2-c)} dx = \sqrt{\frac{\pi}{\delta}} \quad \forall c \in \mathbb{C}, \delta > 0$

$(e^x)^2 = e^{2x}$

$\int_0^1 x \cdot e^{-x^2} = \left[-\frac{1}{2} e^{-x^2} \right]_0^1$

d) (To be discussed in class (2 points)) The average position of a particle at time t is defined as

$\langle x \rangle_t = \int_{-\infty}^{\infty} x |\psi(x, t)|^2 dx$

Show that

$\langle x \rangle_t = v_0 t$

$\langle x \rangle_t = \int_{-\infty}^{\infty} x \cdot |\Psi(x, t)|^2 dx = \frac{1}{a\sqrt{\pi}(1+it/\tau)} \int_{-\infty}^{\infty} x \cdot e^{-\frac{(x-v_0 t)^2}{a^2(1+it/\tau)} + 2i(k_0 x - \omega_0 t)} dx =$

W'keit, dass Teil her...
 $\int_a^b |\Psi(x, t)|^2 dx =$